The spherical MHD code MagIC

Fundamentals

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MHD problem

Installing and running the code

Postprocessing

Outline

1 Introduction

- What for? How?
- Introducting MagIC

2 MHD problem

- 3 Installing and running the code
- 4 Postprocessing

What for?

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Earth's mantle



Solar convective zone

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Earth's core



Cassini

Jupiter

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Glatzmaier & Olson (2005)



Earth's core

Jupiter

Spherical geometry is more natural for studying rotating convection in astrophysical and geophysical objects!

The setup

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Local methods = finite differences, volume, elements?

- PROS: easier to implement, more straightforward to parallelise, grid refinements possible
- CONS: anisotropic grids, pole instability, problem with vacuum magnetic boundary condition, more points required to get same accuracy

Spectral methods =expansion as complete sets of functions?

- PROS: derivatives easier to calculate with high accuracy, magnetic boundary condition is straightforward, lower number of grid points required
- **CONS**: parallelisation harder to implement and more communications

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To date spectral methods are more suitable!

"Local methods [...] need longer elapsed times than spectral methods to achieve the same accuracy with the same number of processors. Spherical harmonic expansion methods [...] offer the best assurance of efficiency for geodynamo simulations" (Matsui et al. 2016)

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Some milestones...

Chandrasekhar (1960s): poloidal/toroidal decomposition, onset of convection in spherical shells

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- Chandrasekhar (1960s): poloidal/toroidal decomposition, onset of convection in spherical shells
- **2** Orszag (1970s): spectral methods in computational fluid dynamics

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- **4 Glatzmaier & Gilman (1980)**: onset of compressible convection in a spherical shell
- **5** Glatzmaier (1984): pseudo-spectral MHD code in a spherical shell geometry

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Pseudo-spectral? What does it mean?

Pseudo-spectral codes

- The linear terms are expanded as complete sets of functions (e.g. spherical harmonics, Chebyshev polynomials, Fourier functions, ...)
- Nonlinear terms treated in grid space rather than spectral space = numerical transformations between spectral and spatial representations

MagIC heritage

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MagIC in words

- MagIC simulates rotating fluid dynamics in a spherical shell
- It solves for the coupled evolution of Navier-Stokes equation, MHD equation, temperature (or entropy) equation and an equation for chemical composition under both the anelastic and the Boussinesq approximations
- A dimensionless formulation of the equations is assumed
- MagIC is a free software (GPL), written in Fortran
- Post-processing relies on python libraries
- Poloidal/toroidal decomposition is employed
- MagIC uses spherical harmonic decomposition in the angular directions
- Chebyshev polynomials or finite differences are employed in the radial direction
- MagIC uses a mixed implicit/explicit time stepping scheme
- The code relies on a hybrid parallelisation scheme (MPI/OpenMP)

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Stucture of the code



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Website and documentation

Since 2015: MagIC is a hosted on https://github.com/magic-sph/magic

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<> Code ① Issue	(2)) Pull requests (0)	III Projects 0	Insights 👻					
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Cmake	Moved common openmp f	ags to the top					14 days a	ago
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icense license	- merge the python subro	utines into the MPI	version (latest version)				2 years a	ago
python/magic	Merge branch 'master' of	https://github.com/	magic-sph/magic				4 days a	ago
samples	fix unit for Graphic output						15 days a	ago
iin src	fix one compiler warning						3 dave :	ano

Online documentation: https://magic-sph.github.io

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Outline

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- Fully compressible equations
- From fully compressible to anelastic
- Dimensionless anelastic equations

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Equation of motion for a compressible fluid

Continuity equation

$$\frac{\partial \rho}{\partial t} + \boldsymbol{\nabla} \cdot (\rho \boldsymbol{u}) = \boldsymbol{0}$$

Navier Stokes equation:

$$\rho\left(\frac{\partial \boldsymbol{u}}{\partial t} + \boldsymbol{u} \cdot \boldsymbol{\nabla} \boldsymbol{u} + 2\boldsymbol{\Omega} \times \boldsymbol{u}\right) = -\boldsymbol{\nabla} \boldsymbol{p} + \rho \, \boldsymbol{g} + \frac{1}{\mu_0} \left(\boldsymbol{\nabla} \times \boldsymbol{B}\right) \times \boldsymbol{B} + \boldsymbol{\nabla} \cdot \boldsymbol{S}$$

with the rate-of-strain tensor expressed by

$$\mathsf{S}_{ij} = \nu \rho \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} - \frac{2}{3} \delta_{ij} \, \boldsymbol{\nabla} \cdot \boldsymbol{u} \right)$$

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Energy equation for a compressible fluid

$$\rho T\left(\frac{\partial s}{\partial t} + \boldsymbol{u} \cdot \boldsymbol{\nabla} s\right) = \boldsymbol{\nabla} \cdot (\boldsymbol{k}_{T} \boldsymbol{\nabla} T) + \boldsymbol{\Phi}_{\nu} + \lambda \left(\boldsymbol{\nabla} \times B\right)^{2} + \epsilon_{T}$$

with the viscous heating Φ_{ν} expressed by

$$\Phi_{
u} = 2
ho\left[e_{ij}e_{ji} - rac{1}{3}(oldsymbol{
abla}\cdotoldsymbol{u})^2
ight]$$

If in addition to that, compositional changes are also considered another equation for the chemical composition ξ reads

$$\rho\left(\frac{\partial\xi}{\partial t}+\boldsymbol{u}\cdot\boldsymbol{\nabla}\xi\right)=\boldsymbol{\nabla}\cdot(\boldsymbol{k}_{\xi}\boldsymbol{\nabla}\xi)+\epsilon_{\xi}$$

Induction equation

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Non-relativistic Maxwell equations provide

$$rac{\partial oldsymbol{B}}{\partial t} = oldsymbol{
abla} imes (oldsymbol{u} imes oldsymbol{B} - \lambda oldsymbol{
abla} imes oldsymbol{B})$$

with $\nabla \cdot \boldsymbol{B} = 0$ When λ is homogeneous, one simply gets

$$\frac{\partial \boldsymbol{B}}{\partial t} = \boldsymbol{\nabla} \times (\boldsymbol{u} \times \boldsymbol{B}) + \lambda \boldsymbol{\Delta} \boldsymbol{B}$$

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Equation of state

In general:

$$p = f(\rho, T, \xi)$$

or

$$\frac{1}{\rho}\partial\rho = -\alpha\,\partial\,T + \beta\,\partial\rho + \delta\partial\xi$$

where

Thermal expansivity:
$$\alpha = -\frac{1}{\rho} \left(\frac{\partial \rho}{\partial T} \right)_{\xi,\rho}$$

Compressibillity: $\beta = \frac{1}{\rho} \left(\frac{\partial \rho}{\partial \rho} \right)_{\xi,\rho}$
Chemical coefficient: $\delta = \frac{1}{\rho} \left(\frac{\partial \rho}{\partial \xi} \right)_{\rho,\rho}$

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From fully compressible to anelastic and Boussinesq Reference state

MHD equations

MagIC either uses the **anelastic** or the **Boussinesq** approximation of the Navier Stokes equation

Anelastic approximation = small disturbance (prime) around an **adiabatic** reference state (tilde):

$$\epsilon \sim rac{s'}{c_{m{p}}} \sim rac{{\cal T}'}{{ ilde {\cal T}}} \sim rac{
ho'}{{ ilde {
ho}}} \sim rac{{m{p}'}}{{ ilde {m{
ho}}}} \sim rac{{m{\xi}'}}{{ ilde {m{\xi}}}}$$

The reference state is hydrostatic, adiabatic, and non magnetic:

$$\boldsymbol{\nabla} \tilde{\boldsymbol{p}} = \tilde{\rho} \, \boldsymbol{g}; \quad \boldsymbol{\nabla} \tilde{\boldsymbol{T}} = \frac{\alpha \, \tilde{\boldsymbol{T}}}{c_{\boldsymbol{p}}} \boldsymbol{g}; \quad \boldsymbol{\nabla} \tilde{\boldsymbol{\xi}} = 0$$

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Anelastic continuity equation

Using $\rho = \tilde{\rho} + \rho'$ yields

$$\underbrace{\frac{\partial \tilde{\rho}}{\partial t}}_{=0} + \frac{\partial \rho'}{\partial t} + \nabla \cdot (\tilde{\rho} \boldsymbol{u}) + \underbrace{\nabla \cdot (\rho' \boldsymbol{u})}_{\mathcal{O}(\epsilon^2)} = 0$$

Estimate of the ratio

$$rac{\partial
ho' / \partial t}{oldsymbol
abla \cdot (ilde{
ho} oldsymbol u)} \sim rac{
ho'}{ ilde{
ho}} \sim \epsilon$$

The first order anelastic equation thus reads

$$oldsymbol{
abla} \cdot (ilde{
ho}oldsymbol{u}) = 0$$

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Anelastic equations

Navier-Stokes equation:

$$\frac{\partial \boldsymbol{u}}{\partial t} + \boldsymbol{u} \cdot \boldsymbol{\nabla} \boldsymbol{u} + 2\boldsymbol{\Omega} \times \boldsymbol{u} = -\boldsymbol{\nabla} \frac{p'}{\tilde{\rho}} - \frac{\tilde{\alpha} \tilde{T}}{c_{\rho}} \boldsymbol{s}' \boldsymbol{g} + \frac{1}{\mu_0 \tilde{\rho}} (\boldsymbol{\nabla} \times \boldsymbol{B}) \times \boldsymbol{B} + \frac{1}{\tilde{\rho}} \boldsymbol{\nabla} \cdot \boldsymbol{S}$$

Energy equation:

$$\tilde{\rho}\tilde{T}\left(\frac{\partial s'}{\partial t}+\boldsymbol{u}\cdot\boldsymbol{\nabla}s'\right)=\boldsymbol{\nabla}\cdot\left(\boldsymbol{k}_{T}\boldsymbol{\nabla}T'\right)+\boldsymbol{\Phi}_{\nu}+\lambda\left(\boldsymbol{\nabla}\times\boldsymbol{B}\right)^{2}+\epsilon_{T}$$

Induction equation:

$$rac{\partial oldsymbol{B}}{\partial t} = oldsymbol{
abla} imes (oldsymbol{u} imes oldsymbol{B} - \lambda oldsymbol{
abla} imes oldsymbol{B})$$

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Boundary conditions

Mechanical boundary conditions:

Stress-free: $\mathbf{n} \times (\mathbf{S} \cdot \mathbf{n}) = \mathbf{0}$, or no-slip: $\mathbf{u} = \mathbf{0}$, $r = r_i, r_o$

Magnetic boundary conditions:

Vacuum:
$$oldsymbol{\Delta} oldsymbol{B} = oldsymbol{0}, \quad r = r_i, r_o$$

Thermal boundary conditions:

Flux:
$$\frac{\partial T'}{\partial r} = 0$$
, or temperature: $T' = 0$, $r = r_i, r_o$

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A dimensionless formulation of the equations is assumed

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A dimensionless formulation of the anelastic MHD equation

MHD equations

MagIC uses a dimensionless form of the anelastic MHD equations

In MagIC, the viscous diffusion time is assumed to be the **reference timescale** and the spherical shell gap the **reference lengthscale**:

$$[\tilde{\rho}] = \tilde{\rho}(r = r_o); \quad [\tilde{T}] = \tilde{T}(r = r_o); \quad [r] = r_o - r_i;$$
$$[t] = \frac{d^2}{\nu}; \quad [u] = \frac{\nu}{d}; \quad [B] = \sqrt{\mu_0 \lambda \tilde{\rho} \Omega}; \quad [p'] = \tilde{\rho}(r = r_o) \frac{\nu^2}{d^2}$$

This implies that the velocity is expressed in **Reynolds number** unit, and the magnetic field in **Elsasser number** unit.

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Dimensionless anelastic MHD equations

In the case of an ideal gas with homogeneous kinematic viscosity ν , thermal diffusivity κ and magnetic diffusivity λ , one gets:

$$\nabla \cdot (\tilde{\rho}\boldsymbol{u}) = 0$$

$$\nabla \cdot \boldsymbol{B} = 0$$

$$\frac{\partial \boldsymbol{u}}{\partial t} + \boldsymbol{u} \cdot \nabla \boldsymbol{u} + \frac{2}{\boldsymbol{E}} \boldsymbol{e}_{\boldsymbol{z}} \times \boldsymbol{u} = -\nabla \frac{\boldsymbol{p}'}{\tilde{\rho}} + \frac{\boldsymbol{R}_{\boldsymbol{a}}}{\boldsymbol{P}_{\boldsymbol{r}}} \boldsymbol{g}(\boldsymbol{r}) \, \boldsymbol{s}' \, \boldsymbol{e}_{\boldsymbol{r}} + \frac{1}{\tilde{\rho} \, \boldsymbol{E} \, \boldsymbol{P}_{\boldsymbol{m}}} \left(\nabla \times \boldsymbol{B} \right) \times \boldsymbol{B} + \frac{1}{\tilde{\rho}} \nabla \cdot \boldsymbol{S}$$

$$\frac{\partial \boldsymbol{B}}{\partial t} = \nabla \times (\boldsymbol{u} \times \boldsymbol{B}) + \frac{1}{\boldsymbol{P}_{\boldsymbol{m}}} \boldsymbol{\Delta} \boldsymbol{B}$$

$$\tilde{\rho} \tilde{T} \left(\frac{\partial \boldsymbol{s}'}{\partial t} + \boldsymbol{u} \cdot \nabla \boldsymbol{s}' \right) = \frac{1}{\boldsymbol{P}_{\boldsymbol{r}}} \nabla \cdot \left(\tilde{\rho} \nabla T' \right) + \frac{\boldsymbol{D} \boldsymbol{i} \, \boldsymbol{P}_{\boldsymbol{r}}}{\boldsymbol{R}_{\boldsymbol{a}}} \left[\Phi_{\nu} + \frac{1}{\boldsymbol{P} \boldsymbol{m}^{2} \, \boldsymbol{E}} \left(\nabla \times \boldsymbol{B} \right)^{2} \right]$$

N.B. In case of compositional convection, another equation and two additional control parameters are required.

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Dimensionless Boussinesq MHD equations

In the Boussinesq limit, $Di \rightarrow 0$, then

$$\nabla \cdot \boldsymbol{u} = 0$$

$$\nabla \cdot \boldsymbol{B} = 0$$

$$\frac{\partial \boldsymbol{u}}{\partial t} + \boldsymbol{u} \cdot \nabla \boldsymbol{u} + \frac{2}{E} \boldsymbol{e}_{\boldsymbol{z}} \times \boldsymbol{u} = -\nabla \boldsymbol{p}' + \frac{Ra}{Pr} \boldsymbol{g}(r) T' \boldsymbol{e}_{r} + \frac{1}{E Pm} (\nabla \times \boldsymbol{B}) \times \boldsymbol{B} + \Delta \boldsymbol{u}$$

$$\frac{\partial \boldsymbol{B}}{\partial t} = \nabla \times (\boldsymbol{u} \times \boldsymbol{B}) + \frac{1}{Pm} \Delta \boldsymbol{B}$$

$$\frac{\partial T'}{\partial t} + \boldsymbol{u} \cdot \nabla T' = \frac{1}{Pr} \Delta T'$$

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From physical properties to dimensionless numbers

Ekman number:
$$E = \frac{\nu}{\Omega d^2}$$
Rayleigh number: $Ra = \frac{\alpha T_o g_o d^3 \Delta s}{c_p \nu \kappa}$ Prandtl number: $Pr = \frac{\nu}{\kappa}$ Magnetic Prandtl number: $Pm = \frac{\nu}{\lambda}$ Dissipation number: $Di = \frac{\alpha T_o g_o}{c_p}$ Radius ratio: $\eta = \frac{r_i}{r_o}$

N.B. when $Di \rightarrow 0$, the Boussinesq limit is recovered.

MHD problem

The (astro/geo)physical regime

Parameter	Earth's core	Giant planets	Sun
E	10^{-15}	10^{-18}	10^{-15}
Ra	10 ²⁷	10 ³⁰	10^{24}
Pr	0.1	0.1	10^{-6}
Pm	10^{-6}	10^{-7}	10^{-3}
Λ (Lorentz/Coriolis)	1	1	?
Ro_ℓ (Inertia/Coriolis)	10^{-2}	10^{-3}	1
Rm (adv./diff.)	1000	10 ⁵	10^{9}
<i>Re</i> (adv./diff.)	10 ⁹	10 ¹²	10^{12}

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What does it actually implies? Is it possible to reach these parameters with my numerical dynamo model?

MHD problem

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Reynolds number: the range of length-scale

$$Re = rac{u_{rms} d}{
u} = rac{d}{\ell_d}$$
 where $\ell_d = rac{
u}{u_{rms}}$
 $\ell_d = rac{d}{Re}$

MHD problem

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Reynolds number: the range of length-scale

$$Re = rac{u_{rms}\,d}{
u} = rac{d}{\ell_d} \quad ext{where} \quad \ell_d = rac{
u}{u_{rms}}$$
 $\ell_d = rac{d}{Re}$

- In natural objects, $I_d \sim 10^{-9} d$
- In other words, the ratio of the bigger length-scale to the smallest one is 10^9 .
- You might need 10^9 grid points in each direction. This implies $Re_{mesh} = 1$.

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Ekman number: the range of time-scales

$$E = rac{
u}{\Omega \, d^2} = rac{P_{rot}}{ au_
u} \quad ext{where} \quad au_
u = rac{d^2}{
u}$$

 τ_{ν} is the viscous diffusion time, $\textit{P}_{\textit{rot}}$ the rotation period.

$$\tau_{\nu} = \frac{P_{rot}}{E}$$

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Ekman number: the range of time-scales

$$E = rac{
u}{\Omega \, d^2} = rac{P_{rot}}{ au_
u} \quad ext{where} \quad au_
u = rac{d^2}{
u}$$

 τ_{ν} is the viscous diffusion time, ${\it P_{rot}}$ the rotation period.

$$\tau_{\nu} = \frac{P_{rot}}{E}$$

- \blacksquare In natural objects, $\tau_{\nu} \sim 10^{15} \, \textit{P}_{rot}$
- In other words, the ratio of the longest time-scale to the smallest one is 10^{15} !
- You might need 10¹⁵ time steps to model the problem

troduction	MHD problem ○○○○○○○○○○○○○○○		Installing and runnir	Postproces	
ummary	,				
	Parameter	Earth's core	Tractable	Hard limit (2015)	
-	Е	10^{-15}	$\geq 10^{-6}$	10 ⁻⁷	
	Ra	10 ²⁷	$\leq 10^{12}$	10 ¹³	
	Pr	0.1	0.1-10	1	
	Pm	10^{-6}	0.1	$6 imes 10^{-2}$	
	Λ (Lorentz/Coriolis)	1	1	1	
	Ro_ℓ (Inertia/Coriolis)	10^{-2}	$10^{-3} - 10^{-1}$	10^{-1}	
	Rm (adv./diff.)	1000	1000	1000	
	<i>Re</i> (adv./diff.)	10^{9}	100-1000	7000	

Two complementary approaches

C

- In the "tractable" regime: parameter studies are possible
- In the "hard-limit" regime, only one single run is possible

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 - Requirements and compilation
 - Executing MagIC

4 Postprocessing

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Requirements to compile MagIC

Requirements

Mandatory Fortran and C compilers

Suggested git (https://git-scm.com/) to clone the code repository

Suggested CMake (https://cmake.org) to build the code

Suggested MPI library: rather use intelMPI or MPICH for full support for hybrid MPI/OpenMP

Optional LAPACK or MKL

Optional SHTns for spherical harmonics transforms

MHD problem

Installing and running the code $0 \bullet 0 \circ 0 \circ 0 \circ 0$

Postprocessing

Data visualisation and post processing Requirements

Post-processing functions are python based. You need to install the following libraries:

Python libraries required

Mandatory matplotlib (https://matplotlib.org): plotting functions
Mandatory scipy (https://www.scipy.org): scientific libraries
Suggested ipython (https://ipython.org): interactive shell
Optional basemap (https://matplotlib.org/basemap/): additional map
projections

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Get the code and compile it

1 Install requirements:

- \$ module load gcc-6 gfortran-6 libopenmpi cmake git
- \$ module load python27 python-scipy ipython python-matplotlib

2 Clone the code from github

\$ git clone https://github.com/magic-sph/magic.git

3 Set-up the environment variables

```
$ cd magic
```

\$ source sourceme.sh # (or sourceme.csh)

4 Define the Fortran and C compilers

- \$ export FC=mpif90 # replace by your compiler
- \$ export CC=mpicc

5 Create a build directory and compile

- \$ mkdir build; cd build
- \$ cmake \$MAGIC_HOME -DUSE_MPI=yes -DUSE_OMP=no
- \$ make -j

MHD problem

Installing and running the code $\circ\circ\circ\bullet\circ\circ\circ\circ$

Postprocessing





Run MagIC

MHD problem

Installing and running the code $\circ\circ\circ\circ\bullet\circ\circ\circ$

Postprocessing

Run with 8 CPUs:

- \$ export OMP_NUM_THREADS=1
- \$ mpiexec -n 8 magic.exe input.nml

input.nml contains all the input informations required to run the code!

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Postprocessing

&grid

Input namelist (1/3)

n_r_max	=33, !	Radial resolution
n_cheb_max	=31, !	Number of Chebyshev polynomials
n_phi_tot	=192, !	Azimuthal resolution
minc	=1, <i>!</i>	Azimuthal symmetry
/		
&control		
mode	=0,	! Magnetic, non-magnetic,
tag	="test",	! Extension of the output files
n_time_steps	s=40000,	! Number of timesteps
dtmax	=1.0D-4,	! Maximum timestep
runHours	=02,	! Run-time
runMinutes	=00,	
/		

Input namelist (2/3)

&phys_param

ra	=1.1D5 ,	!	Rayleigh numbe	r					
ek	=1.0D-3,	!	Ekman number	Ekman number					
pr	=1.0D0,	!	Prandtl number	•					
prmag	=5.0D0	!	Magnetic Prand	lt	l number				
radratio	=0.35D0,	!	Radius ratio r	2	i/r_o				
ktops	=1,	!	BC: fixed-temp	e	rature at the top				
ktopv	=2,	!	BC: rigid wall	, (at the top				
/									
&start_field									
l_start_fil	e=.false.	,		!	Start from a check point?				
start_file	="checkpo	oi	nt_end.start",	!	Name of the check point				
init_b1	=3,			!	Init. mag. field: dipole				
amp_b1	=1,			!	Amplitude \Lambda=1				
init_s1	=0404,			!	Init. temperature perturbation				
amp_s1	=0.03,			!	Amplitude of the init. pert.				

MHD problem

Installing and running the code $\circ\circ\circ\circ\circ\circ\circ\bullet$

Postprocessing

Input namelist (3/3)

&output_conti	rol				
n_log_step	=50,	!	Output	eve	ery n_log_step
n_graphs	=3,	!	Number	of	graphic files
n_rsts	=1,	!	Number	of	restart files
n_stores	=0,				
n_specs	=1,	!	Number	of	spectra
/					
&mantle					
nRotMa	=0				
/					
&inner_core					
sigma_ratio	=1.d0,	!	Conduc	ctir	ng inner-core
nRotIC	=1,	!	Rotati	ing	inner core
/					

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log.TAG file

$\log. {\tt TAG}$ provides all the important information about the run:

- All parameters and other inputs including default values
- Information on parallelization, run time etc
- Log of important events: important output files, changing time step, ...
- Some important time averaged quantities, measures ...



MHD problem

Installing and running the code

Postprocessing

Plotting time series

 $e_kin.TAG$ is always produced. It contains the time evolution of kinetic energy. To plot it:

Open ipython and load the python modules

ipython --matplotlib=gtk (or ipython --pylab)

- >>> from magic import *
- >>> ts = MagicTs(field='e_kin') # Read e_kin.TAG file in \$PWD
- >>> pdoc MagicTs # Gives you the documentation
- Plot the time evolution of magnetic energy

>>> ts = MagicTs(field='e_mag_oc') # Read e_mag_oc.TAG file in \$PWD

Manipulate the data

>>> print(ts.time, ts.emagoc_pol)

MHD problem

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Postprocessing

Loading and plotting snapshots

G_#.TAG files contain 3-D arrays on the grid:

- Load the G_1.TAG file:
 - >>> from magic import *
 - >>> s = Surf(ivar=1)
- Plot the radial velocity u_r in the equatorial plane:

>>> s.equat(field='vr')

■ Plot the ϕ -averaged azimuthal flow u_{ϕ} :

>>> s.avg(field='vp', cm='seismic', levels=33)

• Plot the radial cut of B_r at $r = 0.75 r_o$:

>>> s.surf(field='Br', r=0.75) # Hammer projection

MHD problem

Installing and running the code

Postprocessing 000●0

Data visualisation and post processing Additional outputs

- Plot spectra kin_spec_1.TAG
 - >>> # Plot kin_spec_1.TAG
 - >>> sp = MagicSpectrum(field='kin', ispec=1)
- Plot the time-averaged radial profile of magnetic energy eMagR.TAG

>>> # Plot eMagR.TAG
>>> r = MagicRadial(field='eMagR')

Documentation

And more...

>>> # Movie files (time evolution of 2D slices)
>>> m = Movie()



MHD problem

Installing and running the code

Postprocessing 0000●

Data visualisation and post processing 3-D visualisation with paraview

Requirements

Install a vtk-friendly software: here paraview but VisIt or mayavi should also work fine.

- Read the graphic file you want to convert >>> from magic import MagicGraph >>> gr = MagicGraph(ivar=1) # Load G_1.TAG
 Convert it to a file format readable by paraview >>> # Produce output.vts >>> Graph2Vtk(gr, filename='output')
- 3 Load output.vts with paraview
 - \$ paraview output.vts