# The spherical MHD code MagIC 

Fundamentals

## Thomas Gastine

Institut de Physique du Globe de Paris
6th July 2017


Maglo

## Outline

[1 Introduction

- What for? How?
- Introducting MagIC


## 2 MHD problem

## 3 Installing and running the code

4 Postprocessing

## What for?



Earth's mantle


Solar convective zone

## What for?



Earth's core


Jupiter

## What for?



Earth's core


Jupiter

Spherical geometry is more natural for studying rotating convection in astrophysical and geophysical objects!

## The setup



Rotating spherical shell
Frame of reference rotating with system rotation $\Omega$

Local methods $=$ finite differences, volume, elements?
■ PROS: easier to implement, more straightforward to parallelise, grid refinements possible

- CONS: anisotropic grids, pole instability, problem with vacuum magnetic boundary condition, more points required to get same accuracy

Spectral methods =expansion as complete sets of functions?

- PROS: derivatives easier to calculate with high accuracy, magnetic boundary condition is straightforward, lower number of grid points required
- CONS: parallelisation harder to implement and more communications

To date spectral methods are more suitable!
"Local methods [...] need longer elapsed times than spectral methods to achieve the same accuracy with the same number of processors. Spherical harmonic expansion methods [...] offer the best assurance of efficiency for geodynamo simulations" (Matsui et al. 2016)

1 Chandrasekhar (1960s): poloidal/toroidal decomposition, onset of convection in spherical shells

1 Chandrasekhar (1960s): poloidal/toroidal decomposition, onset of convection in spherical shells
2 Orszag (1970s): spectral methods in computational fluid dynamics

1 Chandrasekhar (1960s): poloidal/toroidal decomposition, onset of convection in spherical shells
2 Orszag (1970s): spectral methods in computational fluid dynamics
3 Young (1974): finite-amplitude convection in a Boussinesq spherical shell using a fully spectral code (roughly $\ell=m=8$ )

1 Chandrasekhar (1960s): poloidal/toroidal decomposition, onset of convection in spherical shells
2 Orszag (1970s): spectral methods in computational fluid dynamics
3 Young (1974): finite-amplitude convection in a Boussinesq spherical shell using a fully spectral code (roughly $\ell=m=8$ )
4 Glatzmaier \& Gilman (1980): onset of compressible convection in a spherical shell

1 Chandrasekhar (1960s): poloidal/toroidal decomposition, onset of convection in spherical shells
2 Orszag (1970s): spectral methods in computational fluid dynamics
3 Young (1974): finite-amplitude convection in a Boussinesq spherical shell using a fully spectral code (roughly $\ell=m=8$ )
4 Glatzmaier \& Gilman (1980): onset of compressible convection in a spherical shell
5 Glatzmaier (1984): pseudo-spectral MHD code in a spherical shell geometry

## Pseudo-spectral? What does it mean?

Pseudo-spectral codes

- The linear terms are expanded as complete sets of functions (e.g. spherical harmonics, Chebyshev polynomials, Fourier functions, ...)
- Nonlinear terms treated in grid space rather than spectral space $=$ numerical transformations between spectral and spatial representations


## MagIC heritage



- MagIC simulates rotating fluid dynamics in a spherical shell
- It solves for the coupled evolution of Navier-Stokes equation, MHD equation, temperature (or entropy) equation and an equation for chemical composition under both the anelastic and the Boussinesq approximations
- A dimensionless formulation of the equations is assumed
- MagIC is a free software (GPL), written in Fortran
- Post-processing relies on python libraries
- Poloidal/toroidal decomposition is employed
- MagIC uses spherical harmonic decomposition in the angular directions
- Chebyshev polynomials or finite differences are employed in the radial direction
- MagIC uses a mixed implicit/explicit time stepping scheme
- The code relies on a hybrid parallelisation scheme (MPI/OpenMP)


## Stucture of the code



## Website and documentation

■ Since 2015: MagIC is a hosted on https://github.com/magic-sph/magic


■ Online documentation: https://magic-sph.github.io

## Outline

## 1 Introduction

2 MHD problem
■ Fully compressible equations

- From fully compressible to anelastic
- Dimensionless anelastic equations

3 Installing and running the code

4 Postprocessing

- It solves for the coupled evolution of Navier-Stokes equation, MHD equation, temperature (or entropy) equation and an equation for chemical composition under both the anelastic and the Boussinesq approximations

A dimensionless formulation of the equations is assumedMagIC is a free software (GPL), written in FortranPost-nrocessing relies on nython librariesPoloidal/toroidal decomposition is employedMagIC uses spherical harmonic decomposition in the angular directionsChebyshev nolynomials or finite differences are employed in the radial direction

- MaglC uses a mixed implicit/explicit time stepping scheme
$\square$ The code relies on a hybrid parallelisation scheme (MPI/OpenMP)


## Equation of motion for a compressible fluid

Continuity equation

$$
\frac{\partial \rho}{\partial t}+\nabla \cdot(\rho \boldsymbol{u})=0
$$

Navier Stokes equation:

$$
\rho\left(\frac{\partial \boldsymbol{u}}{\partial t}+\boldsymbol{u} \cdot \nabla \boldsymbol{u}+2 \boldsymbol{\Omega} \times \boldsymbol{u}\right)=-\nabla p+\rho \boldsymbol{g}+\frac{1}{\mu_{0}}(\nabla \times \boldsymbol{B}) \times \boldsymbol{B}+\nabla \cdot \mathrm{S}
$$

with the rate-of-strain tensor expressed by

$$
\mathrm{S}_{i j}=\nu \rho\left(\frac{\partial u_{i}}{\partial x_{j}}+\frac{\partial u_{j}}{\partial x_{i}}-\frac{2}{3} \delta_{i j} \nabla \cdot \boldsymbol{u}\right)
$$

$$
\rho T\left(\frac{\partial s}{\partial t}+\boldsymbol{u} \cdot \nabla s\right)=\nabla \cdot\left(k_{T} \nabla T\right)+\Phi_{\nu}+\lambda(\nabla \times B)^{2}+\epsilon_{T}
$$

with the viscous heating $\Phi_{\nu}$ expressed by

$$
\Phi_{\nu}=2 \rho\left[e_{i j} e_{j i}-\frac{1}{3}(\boldsymbol{\nabla} \cdot \boldsymbol{u})^{2}\right]
$$

If in addition to that, compositional changes are also considered another equation for the chemical composition $\xi$ reads

$$
\rho\left(\frac{\partial \xi}{\partial t}+\boldsymbol{u} \cdot \boldsymbol{\nabla} \xi\right)=\boldsymbol{\nabla} \cdot\left(k_{\xi} \boldsymbol{\nabla} \xi\right)+\epsilon_{\xi}
$$

Non-relativistic Maxwell equations provide

$$
\frac{\partial \boldsymbol{B}}{\partial t}=\boldsymbol{\nabla} \times(\boldsymbol{u} \times \boldsymbol{B}-\lambda \boldsymbol{\nabla} \times \boldsymbol{B})
$$

with $\boldsymbol{\nabla} \cdot \boldsymbol{B}=0$
When $\lambda$ is homogeneous, one simply gets

$$
\frac{\partial \boldsymbol{B}}{\partial t}=\boldsymbol{\nabla} \times(\boldsymbol{u} \times \boldsymbol{B})+\lambda \boldsymbol{\Delta} \boldsymbol{B}
$$

## Equation of state

In general:

$$
p=f(\rho, T, \xi)
$$

or

$$
\frac{1}{\rho} \partial \rho=-\alpha \partial T+\beta \partial p+\delta \partial \xi
$$

where
Thermal expansivity: $\alpha=-\frac{1}{\rho}\left(\frac{\partial \rho}{\partial T}\right)_{\xi, p}$
Compressibillity: $\beta=\frac{1}{\rho}\left(\frac{\partial \rho}{\partial p}\right)_{\xi, \rho}$
Chemical coefficient: $\delta=\frac{1}{\rho}\left(\frac{\partial \rho}{\partial \xi}\right)_{p, \rho}$

## MHD equations

MagIC either uses the anelastic or the Boussinesq approximation of the Navier Stokes equation

Anelastic approximation $=$ small disturbance (prime) around an adiabatic reference state (tilde):

$$
\epsilon \sim \frac{s^{\prime}}{c_{p}} \sim \frac{T^{\prime}}{\tilde{T}} \sim \frac{\rho^{\prime}}{\tilde{\rho}} \sim \frac{p^{\prime}}{\tilde{p}} \sim \frac{\xi^{\prime}}{\tilde{\xi}}
$$

The reference state is hydrostatic, adiabatic, and non magnetic:

$$
\nabla \tilde{p}=\tilde{\rho} \boldsymbol{g} ; \quad \nabla \tilde{T}=\frac{\alpha \tilde{T}}{c_{p}} \boldsymbol{g} ; \quad \nabla \tilde{\xi}=0
$$

## Anelastic continuity equation

Using $\rho=\tilde{\rho}+\rho^{\prime}$ yields

$$
\underbrace{\frac{\partial \tilde{\rho}}{\partial t}}_{=0}+\frac{\partial \rho^{\prime}}{\partial t}+\boldsymbol{\nabla} \cdot(\tilde{\rho} \boldsymbol{u})+\underbrace{\boldsymbol{\nabla} \cdot\left(\rho^{\prime} \boldsymbol{u}\right)}_{\mathcal{O}\left(\epsilon^{2}\right)}=0
$$

Estimate of the ratio

$$
\frac{\partial \rho^{\prime} / \partial t}{\nabla \cdot(\tilde{\rho} \boldsymbol{u})} \sim \frac{\rho^{\prime}}{\tilde{\rho}} \sim \epsilon
$$

The first order anelastic equation thus reads

$$
\boldsymbol{\nabla} \cdot(\tilde{\rho} \boldsymbol{u})=0
$$

## Anelastic equations

Navier-Stokes equation:

$$
\frac{\partial \boldsymbol{u}}{\partial t}+\boldsymbol{u} \cdot \nabla \boldsymbol{u}+2 \boldsymbol{\Omega} \times \boldsymbol{u}=-\boldsymbol{\nabla} \frac{p^{\prime}}{\tilde{\rho}}-\frac{\tilde{\alpha} \tilde{T}}{c_{p}} s^{\prime} \boldsymbol{g}+\frac{1}{\mu_{0} \tilde{\rho}}(\nabla \times \boldsymbol{B}) \times \boldsymbol{B}+\frac{1}{\tilde{\rho}} \boldsymbol{\nabla} \cdot \mathrm{~S}
$$

Energy equation:

$$
\tilde{\rho} \tilde{T}\left(\frac{\partial s^{\prime}}{\partial t}+\boldsymbol{u} \cdot \nabla s^{\prime}\right)=\boldsymbol{\nabla} \cdot\left(k_{T} \nabla T^{\prime}\right)+\Phi_{\nu}+\lambda(\nabla \times B)^{2}+\epsilon_{T}
$$

Induction equation:

$$
\frac{\partial \boldsymbol{B}}{\partial t}=\boldsymbol{\nabla} \times(\boldsymbol{u} \times \boldsymbol{B}-\lambda \boldsymbol{\nabla} \times \boldsymbol{B})
$$

## Boundary conditions

■ Mechanical boundary conditions:
Stress-free: $\boldsymbol{n} \times(\mathrm{S} \cdot \boldsymbol{n})=\mathbf{0}, \quad$ or no-slip: $\boldsymbol{u}=\mathbf{0}, \quad r=r_{i}, r_{0}$
■ Magnetic boundary conditions:

$$
\text { Vacuum: } \boldsymbol{\Delta} B=\mathbf{0}, \quad r=r_{i}, r_{o}
$$

■ Thermal boundary conditions:
Flux: $\frac{\partial T^{\prime}}{\partial r}=0, \quad$ or temperature: $T^{\prime}=0, \quad r=r_{i}, r_{o}$
$\square$ MagIC simulates rotating fluid dynamics in a spherical shell

- It solves for the coupled evolution of Navier-Stokes equation. MHD equation, temperature (or entropy) equation and an equation for chemical composition under both the anelastic and the Boussinesq approximations
A dimensionless formulation of the equations is assumed
- MagIC is a free software (GPL), written in Fortran
- Post-processing relies on python libraries
- Poloidal/toroidal decomposition is employec
- MagIC uses spherical harmonic decomposition in the angular directions
- Chebyshev polynomials or finite differences are employed in the radial direction
$\square$ MagIC uses a mixed implicit/explicit time stepping scheme
- The code relies on a hybrid parallelisation scheme (MPI/OpenMP)


## A dimensionless formulation of the anelastic MHD equation

## MHD equations

MagIC uses a dimensionless form of the anelastic MHD equations
In MagIC, the viscous diffusion time is assumed to be the reference timescale and the spherical shell gap the reference lengthscale:

$$
\begin{array}{r}
{[\tilde{\rho}]=\tilde{\rho}\left(r=r_{o}\right) ; \quad[\tilde{T}]=\tilde{T}\left(r=r_{o}\right) ; \quad[r]=r_{o}-r_{i} ;} \\
{[t]=\frac{d^{2}}{\nu} ; \quad[u]=\frac{\nu}{d} ; \quad[B]=\sqrt{\mu_{0} \lambda \tilde{\rho} \Omega} ; \quad\left[p^{\prime}\right]=\tilde{\rho}\left(r=r_{o}\right) \frac{\nu^{2}}{d^{2}}}
\end{array}
$$

This implies that the velocity is expressed in Reynolds number unit, and the magnetic field in Elsasser number unit.

## Dimensionless anelastic MHD equations

In the case of an ideal gas with homogeneous kinematic viscosity $\nu$, thermal diffusivity $\kappa$ and magnetic diffusivity $\lambda$, one gets:

$$
\begin{aligned}
\boldsymbol{\nabla} \cdot(\tilde{\rho} \boldsymbol{u}) & =0 \\
\boldsymbol{\nabla} \cdot \boldsymbol{B} & =0 \\
\frac{\partial \boldsymbol{u}}{\partial t}+\boldsymbol{u} \cdot \boldsymbol{\nabla} \boldsymbol{u}+\frac{2}{E} \boldsymbol{e}_{\boldsymbol{z}} \times \boldsymbol{u} & =-\boldsymbol{\nabla} \frac{p^{\prime}}{\tilde{\rho}}+\frac{R a}{P r} g(r) s^{\prime} \boldsymbol{e}_{\boldsymbol{r}}+\frac{1}{\tilde{\rho} E P m}(\boldsymbol{\nabla} \times \boldsymbol{B}) \times \boldsymbol{B}+\frac{1}{\tilde{\rho}} \boldsymbol{\nabla} \cdot \mathrm{~S} \\
\frac{\partial \boldsymbol{B}}{\partial t} & =\boldsymbol{\nabla} \times(\boldsymbol{u} \times \boldsymbol{B})+\frac{1}{P m} \boldsymbol{\Delta} \boldsymbol{B} \\
\tilde{\rho} \tilde{T}\left(\frac{\partial s^{\prime}}{\partial t}+\boldsymbol{u} \cdot \nabla s^{\prime}\right) & =\frac{1}{P r} \boldsymbol{\nabla} \cdot\left(\tilde{\rho} \boldsymbol{\nabla} T^{\prime}\right)+\frac{D i P r}{R a}\left[\Phi_{\nu}+\frac{1}{P m^{2} E}(\boldsymbol{\nabla} \times B)^{2}\right]
\end{aligned}
$$

N.B. In case of compositional convection, another equation and two additional control parameters are required.

## Dimensionless Boussinesq MHD equations

In the Boussinesq limit, $D i \rightarrow 0$, then

$$
\begin{aligned}
\boldsymbol{\nabla} \cdot \boldsymbol{u} & =0 \\
\boldsymbol{\nabla} \cdot \boldsymbol{B} & =0 \\
\frac{\partial \boldsymbol{u}}{\partial t}+\boldsymbol{u} \cdot \boldsymbol{\nabla} \boldsymbol{u}+\frac{2}{E} \boldsymbol{e}_{\mathbf{z}} \times \boldsymbol{u} & =-\boldsymbol{\nabla} p^{\prime}+\frac{R a}{P r} g(r) T^{\prime} \boldsymbol{e}_{r}+\frac{1}{E P m}(\boldsymbol{\nabla} \times \boldsymbol{B}) \times \boldsymbol{B}+\boldsymbol{\Delta u} \\
\frac{\partial \boldsymbol{B}}{\partial t} & =\boldsymbol{\nabla} \times(\boldsymbol{u} \times \boldsymbol{B})+\frac{1}{P m} \boldsymbol{\Delta} \boldsymbol{B} \\
\frac{\partial T^{\prime}}{\partial t}+\boldsymbol{u} \cdot \boldsymbol{\nabla} T^{\prime} & =\frac{1}{P r} \boldsymbol{\Delta} T^{\prime}
\end{aligned}
$$

## From physical properties to dimensionless numbers

Ekman number: $E=\frac{\nu}{\Omega d^{2}}$
Rayleigh number: $R a=\frac{\alpha T_{o} g_{o} d^{3} \Delta s}{c_{p} \nu \kappa}$
Prandtl number: $\operatorname{Pr}=\frac{\nu}{\kappa}$
Magnetic Prandtl number: $\quad P m=\frac{\nu}{\lambda}$
Dissipation number: $\quad D i=\frac{\alpha T_{o} g_{o}}{c_{p}}$

$$
\text { Radius ratio: } \quad \eta=\frac{r_{i}}{r_{o}}
$$

N.B. when $D i \rightarrow 0$, the Boussinesq limit is recovered.

The (astro/geo)physical regime

| Parameter | Earth's core | Giant planets | Sun |
| :---: | :---: | :---: | :---: |
| $E$ | $10^{-15}$ | $10^{-18}$ | $10^{-15}$ |
| $R a$ | $10^{27}$ | $10^{30}$ | $10^{24}$ |
| $P r$ | 0.1 | 0.1 | $10^{-6}$ |
| $P m$ | $10^{-6}$ | $10^{-7}$ | $10^{-3}$ |
| $\Lambda$ (Lorentz/Coriolis) | 1 | 1 | $?$ |
| $R o_{\ell}$ (Inertia/Coriolis) | $10^{-2}$ | $10^{-3}$ | 1 |
| $R m$ (adv./diff.) | 1000 | $10^{5}$ | $10^{9}$ |
| $R e$ (adv./diff.) | $10^{9}$ | $10^{12}$ | $10^{12}$ |

## The (astro/geo)physical regime

| Parameter | Earth's core | Giant planets | Sun |
| :---: | :---: | :---: | :---: |
| $E$ | $10^{-15}$ | $10^{-18}$ | $10^{-15}$ |
| $R a$ | $10^{27}$ | $10^{30}$ | $10^{24}$ |
| $P r$ | 0.1 | 0.1 | $10^{-6}$ |
| $P m$ | $10^{-6}$ | $10^{-7}$ | $10^{-3}$ |
| $\Lambda$ (Lorentz/Coriolis) | 1 | 1 | $?$ |
| $R o_{\ell}$ (Inertia/Coriolis) | $10^{-2}$ | $10^{-3}$ | 1 |
| $R m$ (adv./diff.) | 1000 | $10^{5}$ | $10^{9}$ |
| $R e$ (adv./diff.) | $10^{9}$ | $10^{12}$ | $10^{12}$ |

What does it actually implies? Is it possible to reach these parameters with my numerical dynamo model?

## Reynolds number: the range of length-scale

$$
\begin{gathered}
\operatorname{Re}=\frac{u_{r m s} d}{\nu}=\frac{d}{\ell_{d}} \quad \text { where } \quad \ell_{d}=\frac{\nu}{u_{r m s}} \\
\ell_{d}=\frac{d}{\operatorname{Re}}
\end{gathered}
$$

## Reynolds number: the range of length-scale

$$
\begin{gathered}
\operatorname{Re}=\frac{u_{r m s} d}{\nu}=\frac{d}{\ell_{d}} \quad \text { where } \quad \ell_{d}=\frac{\nu}{u_{r m s}} \\
\ell_{d}=\frac{d}{\operatorname{Re}}
\end{gathered}
$$

- In natural objects, $I_{d} \sim 10^{-9} d$
- In other words, the ratio of the bigger length-scale to the smallest one is $10^{9}$.
- You might need $10^{9}$ grid points in each direction. This implies $R e_{\text {mesh }}=1$.


## Ekman number: the range of time-scales

$$
E=\frac{\nu}{\Omega d^{2}}=\frac{P_{\text {rot }}}{\tau_{\nu}} \quad \text { where } \quad \tau_{\nu}=\frac{d^{2}}{\nu}
$$

$\tau_{\nu}$ is the viscous diffusion time, $P_{r o t}$ the rotation period.

$$
\tau_{\nu}=\frac{P_{r o t}}{E}
$$

## Ekman number: the range of time-scales

$$
E=\frac{\nu}{\Omega d^{2}}=\frac{P_{r o t}}{\tau_{\nu}} \quad \text { where } \quad \tau_{\nu}=\frac{d^{2}}{\nu}
$$

$\tau_{\nu}$ is the viscous diffusion time, $P_{\text {rot }}$ the rotation period.

$$
\tau_{\nu}=\frac{P_{r o t}}{E}
$$

- In natural objects, $\tau_{\nu} \sim 10^{15} P_{\text {rot }}$
- In other words, the ratio of the longest time-scale to the smallest one is $10^{15}$ !
- You might need $10^{15}$ time steps to model the problem


## Summary

| Parameter | Earth's core | Tractable | Hard limit (2015) |
| :---: | :---: | :---: | :---: |
| $E$ | $10^{-15}$ | $\geq 10^{-6}$ | $10^{-7}$ |
| $R a$ | $10^{27}$ | $\leq 10^{12}$ | $10^{13}$ |
| $P r$ | 0.1 | $0.1-10$ | 1 |
| $P m$ | $10^{-6}$ | 0.1 | $6 \times 10^{-2}$ |
| $\Lambda$ (Lorentz/Coriolis) | 1 | 1 | 1 |
| $R o_{\ell}$ (Inertia/Coriolis) | $10^{-2}$ | $10^{-3}-10^{-1}$ | $10^{-1}$ |
| $R m$ (adv./diff.) | 1000 | 1000 | 1000 |
| $R e$ (adv./diff.) | $10^{9}$ | $100-1000$ | 7000 |

Two complementary approaches

- In the "tractable" regime: parameter studies are possible
- In the "hard-limit" regime, only one single run is possible


## Outline

## 1 Introduction

## 2 MHD problem

3 Installing and running the code

- Requirements and compilation
- Executing MagIC

4 Postprocessing

- MaglC simulates rotating fluid dynamics in a spherical shell
$\square$ It solves for the coupled evolution of Navier-Stokes equation, MHD equation, temperature (or entropy) equation and an equation for chemical composition under both the anelastic and the Boussinesq approximations
- A dimensionless formulation of the equations is assumed

MagIC is a free software (GPL), written in Fortran

- Post-processing relies on python libraries
- Poloidal/toroidal decomposition is employed
- MagIC uses spherical harmonic decomposition in the angular directions
- Chebyshev nolynomials or finite differences are employed in the radial direction
- MaglC uses a mixed implicit/explicit time stepping scheme
$\square$ The code relies on a hybrid parallelisation scheme (MPI/OpenMP)


## Requirements to compile MagIC

## Requirements

Mandatory Fortran and C compilers
Suggested git (https://git-scm.com/) to clone the code repository
Suggested CMake (https://cmake.org) to build the code
Suggested MPI library: rather use intelMPI or MPICH for full support for hybrid MPI/OpenMP
Optional LAPACK or MKL
Optional SHTns for spherical harmonics transforms

Post-processing functions are python based. You need to install the following libraries:
Python libraries required
Mandatory matplotlib (https://matplotlib.org): plotting functions
Mandatory scipy (https://www.scipy.org): scientific libraries
Suggested ipython (https://ipython.org): interactive shell
Optional basemap (https://matplotlib.org/basemap/): additional map projections

## Get the code and compile it

1 Install requirements:
\$ module load gcc-6 gfortran-6 libopenmpi cmake git
\$ module load python27 python-scipy ipython python-matplotlib
$\square$ Clone the code from github
\$ git clone https://github.com/magic-sph/magic.git
3 Set-up the environment variables
\$ cd magic
\$ source sourceme.sh \# (or sourceme.csh)
4 Define the Fortran and C compilers
\$ export $\mathrm{FC}=\mathrm{mpif} 90$ \# replace by your compiler
\$ export $\mathrm{CC}=\mathrm{mpicc}$
5 Create a build directory and compile
\$ mkdir build; cd build
\$ cmake \$MAGIC_HOME -DUSE_MPI=yes -DUSE_OMP=no
\$ make -j

## MagIC structure



## Run MagIC

## Run with 8 CPUs:

\$ export OMP_NUM_THREADS=1
\$ mpiexec -n 8 magic.exe input.nml
input.nml contains all the input informations required to run the code!

## Input namelist (1/3)

```
&grid
    n_r_max =33, ! Radial resolution
    n_cheb_max =31, ! Number of Chebyshev polynomials
    n_phi_tot =192, ! Azimuthal resolution
    minc =1, ! Azimuthal symmetry
/
&control
    mode =0, ! Magnetic, non-magnetic, ...
    tag ="test", ! Extension of the output files
    n_time_steps=40000, ! Number of timesteps
    dtmax =1.0D-4, ! Maximum timestep
    runHours =02, ! Run-time
    runMinutes =00,
/
```


## Input namelist (2/3)

\&phys_param

$$
\begin{array}{lll}
\text { ra } & =1.1 \mathrm{D} 5, \text { ! Rayleigh number } \\
\text { ek } & =1.0 \mathrm{D}-3, \text { Ekman number } \\
\text { pr } & =1.0 \mathrm{DO}, \text { ! Prandtl number } \\
\text { prmag } & =5.0 \mathrm{D} 0 \text { ! Magnetic Prandtl number } \\
\text { radratio } & =0.35 \mathrm{DO}, \text { ! Radius ratio } r_{-} i / r_{-} \\
\text {ktops } & =1, & \text { BC: fixed-temperature at the top } \\
\text { ktopv } & =2, \quad \text { BC: rigid wall at the top }
\end{array}
$$

/
\&start_field

```
l_start_file=.false., ! Start from a check point?
start_file ="checkpoint_end.start", ! Name of the check point
init_b1 =3, ! Init. mag. field: dipole
amp_b1 =1,
init_s1 =0404, ! Init. temperature perturbation
amp_s1 =0.03, ! Amplitude of the init. pert.
```


## Input namelist (3/3)

```
&output_control
    n_log_step =50, ! Output every n_log_step
    n_graphs =3, ! Number of graphic files
    n_rsts =1, ! Number of restart files
    n_stores =0,
    n_specs =1, ! Number of spectra
&mantle
    nRotMa =0
/
&inner_core
    sigma_ratio =1.d0, ! Conducting inner-core
    nRotIC =1, ! Rotating inner core
/
```


## Outline

## 1 Introduction

## 2 MHD problem

3 Installing and running the code

4 Postprocessing
log.TAG provides all the important information about the run:

- All parameters and other inputs including default values
- Information on parallelization, run time etc

■ Log of important events: important output files, changing time step, ...

- Some important time averaged quantities, measures ...
e_kin.TAG is always produced. It contains the time evolution of kinetic energy. To plot it:

■ Open ipython and load the python modules

```
ipython --matplotlib=gtk (or ipython --pylab)
>>> from magic import *
>>> ts = MagicTs(field='e_kin') # Read e_kin.TAG file in $PWD
>>> pdoc MagicTs # Gives you the documentation
```

- Plot the time evolution of magnetic energy

```
>>> ts = MagicTs(field='e_mag_oc') # Read e_mag_oc.TAG file in $PWD
```

- Manipulate the data

```
>>> print(ts.time, ts.emagoc_pol)
```


## Loading and plotting snapshots

## G_\#.TAG files contain 3-D arrays on the grid:

- Load the G_1.TAG file:

```
>>> from magic import *
>>> s = Surf(ivar=1)
```

■ Plot the radial velocity $u_{r}$ in the equatorial plane:

```
>>> s.equat(field='vr')
```

■ Plot the $\phi$-averaged azimuthal flow $u_{\phi}$ :

```
>>> s.avg(field='vp', cm='seismic', levels=33)
```

- Plot the radial cut of $B_{r}$ at $r=0.75 r_{0}$ :

```
>>> s.surf(field='Br', r=0.75) # Hammer projection
```


## Data visualisation and post processing

■ Plot spectra kin_spec_1.TAG

```
>>> # Plot kin_spec_1.TAG
>>> sp = MagicSpectrum(field='kin', ispec=1)
```

- Plot the time-averaged radial profile of magnetic energy eMagR.TAG

```
>>> # Plot eMagR.TAG
>>> r = MagicRadial(field='eMagR')
```

- And more...

```
>>> # Movie files (time evolution of 2D slices)
>>> m = Movie()
```


## Requirements

Install a vtk-friendly software: here paraview but VisIt or mayavi should also work fine.

1 Read the graphic file you want to convert
>>> from magic import MagicGraph
>>> gr = MagicGraph(ivar=1) \# Load G_1.TAG
2 Convert it to a file format readable by paraview
>>> \# Produce output.vts
>>> Graph2Vtk(gr, filename='output')
3 Load output.vts with paraview
\$ paraview output.vts

